

3.5

Consider a simply connected volume,  $\mathcal{V}$ , with charge density  $\rho$ . At each boundary of  $\mathcal{V}$ , either  $V$  or  $\frac{\partial V}{\partial n} = \vec{\nabla} V \cdot \hat{n}$  is given. Then  $\exists$  a unique solution to Poisson's equation in this region.

Proof: Let  $V_1$  &  $V_2$  denote two solutions to  $\nabla^2 V = \frac{1}{\epsilon_0} \rho$  in the region specified.

$$\text{Then } \nabla^2 V_1 = \frac{1}{\epsilon_0} \rho \quad \text{and} \quad \nabla^2 V_2 = \frac{1}{\epsilon_0} \rho$$

$$\text{define } V_3 \equiv V_1 - V_2. \quad \text{Then } \nabla^2 V_3 = \nabla^2 V_1 - \nabla^2 V_2 = \frac{1}{\epsilon_0} \rho - \frac{1}{\epsilon_0} \rho = 0.$$

So  $V_3$  satisfies not Poisson's eq., but Laplace's equation.

At every boundary, either  $V_1 = V_2$  and so  $V_3 = 0$

$$\text{or } \frac{\partial V_1}{\partial n} = \frac{\partial V_2}{\partial n} \quad \text{meaning} \quad \frac{\partial V_3}{\partial n} = \frac{\partial}{\partial n} (V_1 - V_2) = \frac{\partial V_1}{\partial n} - \frac{\partial V_2}{\partial n} = 0.$$

Therefore, at every boundary  $V_3$  is zero, or,  $\frac{\partial V_3}{\partial n}$  is 0.

The case  $V_3 = 0$  has been shown in the book. For  $\frac{\partial V_3}{\partial n} = 0$ :

$$\vec{\nabla} \cdot (V_3 \vec{\nabla} V_3) = V_3 \nabla^2 V_3 + \vec{\nabla} V_3 \cdot (\vec{\nabla} V_3)$$

and since  $\nabla^2 V_3 = 0$  this leaves

$$\vec{\nabla} \cdot (V_3 \vec{\nabla} V_3) = E_3^2. \quad \text{Taking a closed volume integral of both sides,}$$

$$\oint_{\mathcal{V}} \vec{\nabla} \cdot (V_3 \vec{\nabla} V_3) d\tau = \oint_{\mathcal{V}} E_3^2 d\tau \Rightarrow \oint_{\mathcal{V}} V_3 (\vec{\nabla} V_3) \cdot d\vec{a} = \oint_{\mathcal{V}} E_3^2 d\tau$$

$$\text{but } (\vec{\nabla} V_3) \cdot d\vec{a} = (\vec{\nabla} V_3) \cdot \hat{n} da \quad \text{and} \quad \vec{\nabla} V_3 \cdot \hat{n} \equiv \frac{\partial V_3}{\partial n} = 0 \quad \text{so}$$

$$\oint_{\mathcal{V}} E_3^2 d\tau = 0. \quad \text{This is only possible if } E_3 = 0.$$

$$\text{With } V_3 = V_1 - V_2 \quad \vec{\nabla} V_3 = \vec{\nabla} V_1 - \vec{\nabla} V_2 \Rightarrow \vec{E}_3 = \vec{E}_2 - \vec{E}_1$$

And hence  $\vec{E}_2 = \vec{E}_1$ , a unique  $\vec{E}$  guarantees a unique  $V$ . Since  $V_1 = \int \vec{E}_1 \cdot d\vec{l} = \int \vec{E}_2 \cdot d\vec{l} = V_2$ .  $\square$